

(3)  $\Delta^{++}$ : uuu All  $\uparrow\uparrow\uparrow$  configuration

→ Pauli exclusion principle violated?

→ solution: color label (RGB)

$$\Psi_{\Delta} = \Psi_{\text{flavor}} \cdot \Psi_{\text{spin}} \cdot \Psi_{\text{space}} \cdot \Psi_{\text{color}}$$

$\downarrow$                      $\downarrow$                      $\downarrow$                      $\downarrow$   
 even                even                even                hence odd!

→ antisym.

RGB - GRB + GBR - BGR  
etc.

2. Use Lorentz Invariance:

$$E_{\text{lab, tot}}^2 - P_{\text{lab, tot}}^2 = E_{\text{cm, tot}}^2 - P_{\text{cm, tot}}^2 = M_{\Delta}^2$$

$$(E_{\pi} + M_p)^2 - P_{\pi}^2 = M_{\Delta}^2$$

$$E_{\pi}^2 + M_p^2 + 2M_p E_{\pi} - P_{\pi}^2 = M_{\Delta}^2$$

$$M_{\pi}^2 + M_p^2 + 2M_p E_{\pi} = M_{\Delta}^2 \Rightarrow E_{\pi} = \frac{M_{\Delta}^2 - M_{\pi}^2 - M_p^2}{2M_p}$$

$$= 329.8 \text{ MeV}$$

$$\Rightarrow P_{\pi} = \sqrt{E_{\pi}^2 - M_{\pi}^2} = \underline{\underline{299 \text{ MeV}/c}}$$

(3) cont'd

(2)

$$3. \quad \Delta \xrightarrow{L} \bar{u} + N$$

$$I(J^P): \frac{3}{2}(\frac{3}{2}^+) \rightarrow 1(0^-) + \frac{1}{2}(\frac{1}{2}^+)$$

$$P_{\text{rel}} = P_{\bar{u}} \cdot P_N \cdot (-1)^L = (-1)^{L+1} \Rightarrow L \text{ odd}$$

$$\frac{J=L+1}{L=1} : 1 + \frac{1}{2} = (\frac{1}{2}, \frac{3}{2}) \text{ o.k.}$$

$$L=3 : 3 + \frac{1}{2} = (\frac{5}{2}, \frac{7}{2}) \Rightarrow N_0$$

etc

L=1 possible  
P-wave.

$$4. \quad \Delta^0 \rightarrow \bar{u}^0 + n$$

$$\rightarrow \bar{u} + p$$

Evaluate isospin multiplets:

$$\begin{array}{c}
 \begin{array}{c} \bar{\pi} \\ (I=1) \end{array} \begin{array}{c} \bar{u} \quad \pi^0 \quad u \\ -1 \quad 0 \quad +1 \end{array} \rightarrow I_3 \\
 \\
 \begin{array}{c} N \\ (I=\frac{1}{2}) \end{array} \begin{array}{c} n \quad p \\ -\frac{1}{2} \quad +\frac{1}{2} \end{array} \rightarrow I_3 \\
 \\
 \begin{array}{c} (I=\frac{3}{2}) \end{array} \begin{array}{c} (\bar{\pi}n) \quad (\bar{u}^0n) \quad (\bar{u}^0p) \quad (\bar{u}^+p) \\ -\frac{3}{2} \quad -\frac{1}{2} \quad \frac{1}{2} \quad \frac{3}{2} \end{array} \rightarrow I_3 \\
 \\
 \begin{array}{c} (I=\frac{1}{2}) \end{array} \begin{array}{c} (\bar{u}p) \quad (\bar{u}^0p) \\ -\frac{1}{2} \quad \frac{1}{2} \end{array} \rightarrow I_3
 \end{array}$$

couples to  $\Delta^0$

$$T_+ |\bar{\pi}n\rangle = (T_{+\bar{u}} + T_{+n}) |\bar{\pi}n\rangle$$

$$= [\sqrt{2} |\bar{u}^0n\rangle + 1 \cdot |\bar{\pi}p\rangle] / \sqrt{3}$$

normalized

$$\left. \begin{array}{l}
 \langle \Delta^0 | \bar{u}^0n \rangle^2 = \frac{2}{3} \\
 \langle \Delta^0 | \bar{u}p \rangle^2 = \frac{1}{3}
 \end{array} \right\} \text{Branching fractions } B$$

(3) cont'd

(3)

$$\sigma = \left( \frac{2\pi}{p^2} \right) \frac{\Gamma_{\Delta \rightarrow \bar{u}p} \Gamma_{\Delta \rightarrow \bar{u}n}}{(E_{cm} - M_\Delta)^2 + \Gamma^2/4}$$

Note  $\Gamma_{\Delta \rightarrow \bar{u}p} = \mathcal{B}_{\Delta \rightarrow \bar{u}p} \cdot \Gamma = \frac{1}{3} \Gamma$

$$\Gamma_{\Delta \rightarrow \bar{u}n} = \mathcal{B}_{\Delta \rightarrow \bar{u}n} \Gamma = \frac{2}{3} \Gamma$$

$$E_{cm} = M_\Delta$$

$$\Rightarrow \sigma = \left( \frac{2\pi}{p^2} \right) \frac{\frac{2}{9} \Gamma^2}{\Gamma^2/4} = \frac{16\pi}{9p^2} \otimes (hc)^2$$

↳ to get right units (barns)

$p^?$

$\xrightarrow{\bar{u}^-} \xleftarrow{p}$  in C.M. frame:

(1)  $E_{\bar{u}} + E_p = M_\Delta$  (E-conservation)

(2)  $|P_{\bar{u}}| = |P_p| \equiv P$  (mom.-conservation)

(2)  $P_{\bar{u}}^2 = P_p^2 \Leftrightarrow P_{\bar{u}}^2 + m_{\bar{u}}^2 + m_p^2 = P_p^2 + m_p^2 + m_{\bar{u}}^2$

$\Leftrightarrow E_{\bar{u}}^2 + m_p^2 = E_p^2 + m_{\bar{u}}^2$

+ (1)  
( $\Rightarrow$ )  $E_{\bar{u}}^2 + m_p^2 = (M_\Delta - E_{\bar{u}})^2 + m_{\bar{u}}^2$

$$= M_\Delta^2 + E_{\bar{u}}^2 - 2M_\Delta E_{\bar{u}} + m_{\bar{u}}^2$$

$$\Rightarrow E_{\bar{u}} = \frac{M_\Delta^2 + m_{\bar{u}}^2 - m_p^2}{2M_\Delta} = \frac{2667}{2277} \text{ MeV}$$

$$\Rightarrow P_{\bar{u}} = P_p = p = \sqrt{E_{\bar{u}}^2 - m_{\bar{u}}^2} = \frac{197}{227.7} \text{ MeV}$$

(1b = 100 fm<sup>2</sup>)

$$\Rightarrow \sigma = \frac{16\pi}{9(227.7)^2 \text{ MeV}^2} \cdot (197)^2 \text{ MeV}^2 \text{ fm}^2 = 4.2 \text{ fm}^2 = 4.2 \cdot 10^{-2} \text{ barns}$$

$$= \underline{\underline{42 \text{ mb}}}$$

IIa

$$E_{cm} = \sqrt{s} = 270 + 270 = 540 \text{ GeV} = E_{jet} + E_W$$

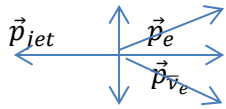
$$\vec{p}_{jet} + \vec{p}_W = 0$$

$$E_{cm} = E_{jet} + E_W = p_W + \sqrt{p_W^2 + m_W^2} \rightarrow (E_{cm} - p_W)^2 = p_W^2 + m_W^2$$

$$p_W = \frac{E_{cm}^2 - m_W^2}{2E_{cm}}$$

$$\vec{p}_W = \vec{p}_e + \vec{p}_{\bar{\nu}_e}$$

Energy and momentum of electron and neutrino must be equal for maximum opening angle.



Use invariant mass to get the opening angle

$$m_W^2 = (E_e + E_{\nu})^2 - (\vec{p}_e + \vec{p}_{\nu})^2$$

$$|p_e| = |p_{\nu}| = E_e = E_{\nu} = \frac{E_W}{2} = \frac{E_{cm} - |p_W|}{2} = \frac{E_{cm} - \frac{E_{cm}^2 - m_W^2}{2E_{cm}}}{2} = \frac{E_{cm}^2 + m_W^2}{4E_{cm}} = x$$

$$m_W^2 = (2x)^2 - (2x^2 + 2x^2 \cos \theta_{e\nu}) = 2x^2(1 - \cos \theta_{e\nu})$$

$$\theta_{e\nu} = \text{acos} \left( 1 - \frac{m_W^2}{2x^2} \right) = \text{acos} \left( 1 - \frac{8E_{cm}^2 m_W^2}{(E_{cm}^2 + m_W^2)^2} \right) = \text{acos} \left( 1 - \frac{8 * 540^2 * 80^2}{(540^2 + 80^2)^2} \right) = 34^\circ$$

IIb (for answers on b and c see pages 230 and 231 of book)

The fact that  $W \rightarrow l^+ \nu_l$  has the same branching for all  $l = e, \mu, \tau \approx 10.8\%$

IIc

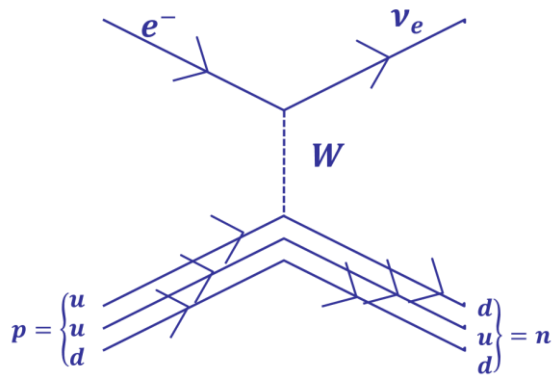
$W^+ \rightarrow u\bar{d}'$  or  $c\bar{s}'$  and not  $t\bar{b}'$  ( $m_t > m_W$ ) each quark pair has three possibilities (color) means 6 possibilities to make hadrons as compared to a lepton pair  $6 * 10.8\% = 65\%$  is close to the number in the PDG datasheet (68% , we ignore all details of phase space and higher order diagrams)

The Cabbibo angle does not play a role the generation does not enter in describing the  $d'$  and  $s'$  thus we working with an approximate unitary transformation going from  $(d, s) \rightarrow (d', s')$ .

IId (see book p 223)

$$\Gamma_{\nu_l} = \alpha_W M_W = 10.8\% \text{ of } 2.08 \text{ GeV} \quad \alpha_W = \frac{0.108 * 2.08}{80} = 2.79 \times 10^{-3} \text{ or } \frac{1}{\alpha_w} = 360$$

IIIa



Determine energy threshold:

$$E_e + m_p = E_\nu + E_n \xrightarrow{\text{threshold}} E_n \approx m_n + \frac{p_e^2}{2m_n}$$

$$\vec{p}_e = \vec{p}_\nu + \vec{p}_n \xrightarrow{\text{threshold}} \vec{p}_n$$

Ignoring recoil threshold is  $E_e = m_n - m_p = 2.933$  MeV with recoil solve quadratic equation

$$E_e + m_p = m_n + \frac{E_e^2 - m_e^2}{2m_n}$$

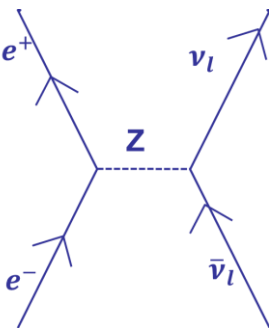
$$E_e = m_n - m_p - \frac{m_e^2}{2m_n} = 2.932 \text{ MeV}$$

Electron neutrino is born and propagates as mass eigenstates. May assume independent propagations (incoherent) therefore  $0.85^2 m_1$  and  $0.53^2 m_2$  and  $0 m_3$  note that total probability is 1 (one). When they arrive on Earth

$m_1$  has probability  $0.85^2$  to be an electron neutrino,  $m_2$  has  $0.53^2$  electron probability

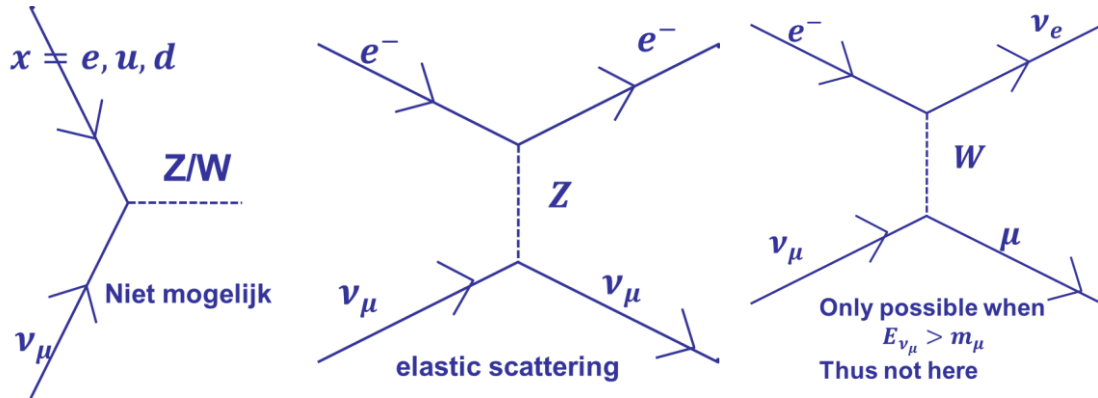
Thus we have a probability  $0.85^4 + 0.53^4 = 0.60$  to be an electron neutrino.

IIIb



Equal amounts of each neutrino and antineutrino i.e.  $1/6$  after propagation because of unitarity also  $1/6$

IIIc



Only middle diagram possible.

III d

$$\Delta t = \frac{L}{v_1} - \frac{L}{v_2} = \frac{L}{c} \left( \frac{1}{\beta_1} - \frac{1}{\beta_2} \right)$$

$$\beta = \frac{p}{E} = \frac{\sqrt{E^2 - m^2}}{E} \approx 1 - \frac{1}{2} \frac{m^2}{E^2}$$

$$\begin{aligned} \Delta t &= \frac{L}{c} \left( 1 + \frac{1}{2} \frac{m^2}{E_1^2} - 1 - \frac{1}{2} \frac{m^2}{E_2^2} \right) \rightarrow \frac{m^2 L}{2c} \left( \frac{1}{E_1^2} - \frac{1}{E_2^2} \right) \rightarrow m^2 = \frac{\frac{2c\Delta t}{L}}{\left( \frac{1}{E_1^2} - \frac{1}{E_2^2} \right)} \\ &= 2 \times \frac{5.3}{5.3} \times 10^{-12} \times \left( 1 - \frac{1}{9} \right) [MeV]^2 \rightarrow m = \sqrt{\frac{18}{8}} \approx 1.4 \text{ eV} \end{aligned}$$

III e p330

$\langle m \rangle = \sum_i |U_{ei}|^2 m_i$  (the weighted sum (sum of weights is one) of probabilities that  $m_i$  contributes)

III d

$$\langle m^2 \rangle = \sum_i |U_{ei}|^2 m_i^2$$

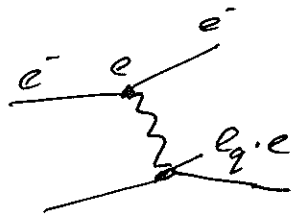
Same argument but for  $m_i^2$ .

(4)

(4)

1.  $\delta$ -function: classic scattering of electron on particle with mass  $m \Rightarrow x = \frac{Q^2}{2m_0} = 1$  (E2p conservation)  
 Hence,  $\delta$  function guarantees that  $\underline{x=1}$ . ~~oh~~ Note,  $Q^2 = M^2 + 2M_0 - W^2$   
 For elastic scattering,  $M=W \Rightarrow Q^2 = 2M_0$

$e_i^2$ ; follows from  $\alpha(n e^2) \rightarrow e_i \alpha$ , since:



2. For a proton  $i$ :

$$W_2^i = e_i^2 \delta\left(\nu - \frac{Q^2}{2EM_p}\right)$$

$\Rightarrow$  incoherent sum with probability density  $f_i(z)$ :

$$\begin{aligned} W_2 &= \sum_i \int e_i^2 \delta\left(\nu - \frac{Q^2}{2EM_p}\right) f_i(z) dz \\ &= \delta\left[\frac{\nu}{z} \left(z - \frac{Q^2}{2M_0}\right)\right] \\ &= \delta\left[\frac{\nu}{z} (z-x)\right] \\ &= \frac{z}{\nu} \delta(z-x) \end{aligned}$$

$$\begin{aligned} W_2 &= \sum_i \int e_i^2 \frac{z}{\nu} \delta(z-x) f_i(z) dz = \frac{1}{\nu} \sum_i e_i^2 \int z f_i(z) \delta(z-x) dz \\ &= \frac{1}{\nu} \sum_i e_i^2 x f_i(x) \end{aligned}$$

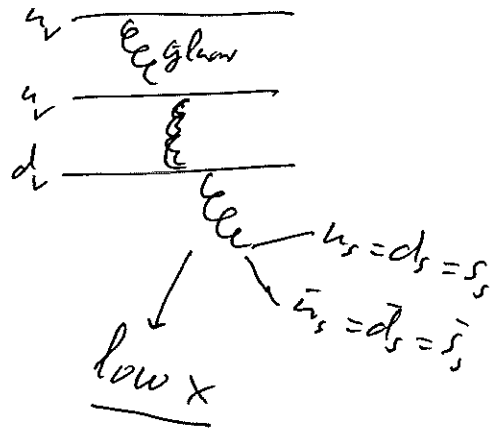
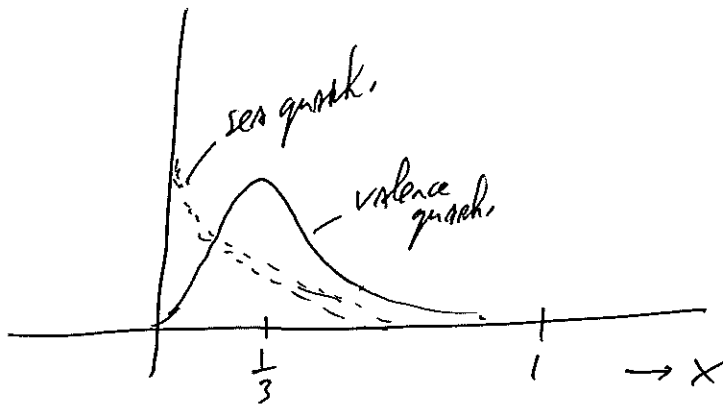
$$\Rightarrow \nu \cdot W_2 = \sum_i e_i^2 x f_i(x) \equiv F_2(x)$$

(4) cont'd

(5)

3. The ~~structure~~ <sup>form factor</sup> function  $F_2$  (also  $F_1 = M_p W_1$ ) only depends upon 1 variable,  $x$ , independent on  $Q^2$ !  
 $\Rightarrow$  parton is point-like

4.



valence quark  
around  $\frac{1}{3}$ , but smeared due to gluon exchange.